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## Study of the structure, magnetic properties and free energy of the three phase composites $La_{0.7-v-z}Sr_{0.3-x}Mn_{1-\delta/3}O_{3-\delta}/(La_2O_3)_{v/2}/(La(OH)_3)_z$

S.P. Liu, G.D. Tang\*, Z.Z. Li, W.H. Qi, D.H. Ji, Y.F. Li, W. Chen, D.L. Hou

Hebei Advanced Thin Films Laboratory, Department of Physics, Hebei Normal University, Shijiazhuang City 050016, People's Republic of China

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#### ABSTRACT

The structural and magnetic properties and the free energy of manganites with the nominal composition La<sub>0.7</sub>Sr<sub>0.3-x</sub>MnO<sub>3- $\delta$ </sub> ( $0 \le x \le 0.10$ ) have been investigated. X-ray diffraction patterns indicate that the materials possess three phases with the  $R\overline{3}c$  perovskite being the dominant phase, La<sub>2</sub>O<sub>3</sub> being the second phase and La(OH)3 being the third phase. Due to the samples were prepared using a solid-state reaction method with the highest heat treatment temperature being 1173 K, we assume that there were no vacancies in the perovskite phase. These materials can be represented as  $La_{0.7-y-z}Sr_{0.3-x}Mn_{1-\delta/3}O_{3-\delta}/(La_2O_3)_{v/z}/(La(OH)_3)_z$ , with  $\delta = 3(x+y+z)$ . Magnetization measurements showed that the dependence of the Curie temperature  $T_C$  on the Mn<sup>4+</sup> ion content, is similar to that of the typical perovskite manganite  $La_{1-x}Sr_xMnO_3$ . The paramagnetic-ferromagnetic transition region of the samples is wider than 315 K. We have also calculated the free energy of the perovskite phase, and explained successfully the dependences on the doping level x of the unit cell volume in the perovskite phase.

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#### 1. Introduction

Colossal magnetoresistance (CMR) materials [1–4] with the general formula  $Ln_{1-x}M_xMnO_3$ , where Ln is the lanthanide cation and M is usually an alkaline-earth cation, have attracted much attention because of their potential applications in magnetic sensors, data storage, and other applications, such as the cathodes of the solid oxide fuel cells [5–7] (SOFCs) and microwave dielectric materials [8–10]. These materials possess an ABO<sub>3</sub> perovskite structure. Zener proposed that the spin structure and the electronic properties of these compounds are correlated via the double exchange (DE) mechanism. According to this picture, in the configuration Mn<sup>3+</sup>-O<sup>2-</sup>-Mn<sup>4+</sup>, the easy simultaneous transfer of an electron (DE carrier) from  $Mn^{3+}$  to  $O^{2-}$  to  $Mn^{4+}$ , causes at the same time a high electrical conductivity and, by the tendency of the traveling electron to retain its spin orientation, a parallel orientation of the magnetic moments of the Mn<sup>3+</sup> and Mn<sup>4+</sup> ions [11,1,2].

Previous study in our group has demonstrated methods of phase analysis and calculating the various ion contents of perovskite phase in lanthanum deficiency in La<sub>0.7-x</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> composition samples with Mn<sub>3</sub>O<sub>4</sub> second phase [12]. In this paper, we investigate the self-doped perovskite manganites with nominal composition La<sub>0.7</sub>Sr<sub>0.3-x</sub>MnO<sub>3</sub> prepared by a solid-state reaction method with the highest heat treatment temperature being 1173 K It is found that these materials can be represented as  $La_{0.7-v-z}Sr_{0.3-x}Mn_{1-\delta/3}O_{3-\delta}/(La_2O_3)_{v/2}/(La(OH)_3)_z$ , with  $\delta = 3(x+y+z)$ .

#### 2. Experimental details

Polycrystalline samples with the nominal composition  $La_{0.7}Sr_{0.3-x}MnO_3$  (x = 0, 0.05, 0.075, and 0.1) were synthesized using the solid-state reaction method. The starting materials were La<sub>2</sub>O<sub>3</sub>, SrCO<sub>3</sub>, and MnO<sub>2</sub> powders. The La<sub>2</sub>O<sub>3</sub> was fired in air at 1073 K for 3 h before use in order to ensure that is dry. The starting materials were intimately mixed in an agate mortar, grinding for 3 h, and heated in air at temperatures up to 1173 K for 24 h, followed by furnace cooling. The samples studied in this work were obtained after regrinding.

Phase identification of the samples was performed by X-ray diffraction (XRD) with an X'pert Pro diffractometer using Cu  $K_{\alpha}$  radiation. The data were collected in the  $2\theta$  range 20– $120^{\circ}$  with a step size of  $0.0167^{\circ}$ . The magnetization measurements were carried out using a Quantum Design Physical Property Measurement System (PPMS), in the temperature range 5-400 K. The temperature at which the dM/dTcurve tends to zero with increasing temperature indicates the Curie temperature  $(T_{\mathbb{C}})$ . The Curie temperature measurements were performed at an applied field of 0.05 T. In order to facilitate the discussion, we label the samples with nominal compositions  $La_{0.7}Sr_{0.3-x}MnO_3$  (x = 0.00, 0.05, 0.075, and 0.10) as N1, N2, N3, and N4, respectively.

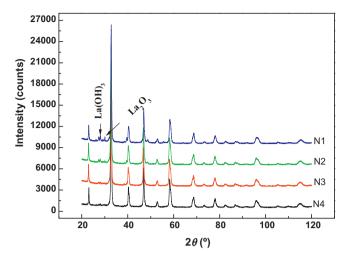
#### 3. Results and discussion

### 3.1. X-ray diffraction analyses

Fig. 1 shows the X-ray diffraction patterns for the samples. It can be seen that all the samples contain three phases with the  $R\overline{3}c$ 

E-mail address: tanggd@mail.hebtu.edu.cn (G.D. Tang).

<sup>\*</sup> Corresponding author.



**Fig. 1.** XRD diffraction patterns of the samples with nominal composition  $La_{0.7}Sr_{0.3-x}MnO_{3-\delta}$ . The main diffraction peaks of the  $La_2O_3$  phase and  $La(OH)_3$  phase are marked by the symbol " $\downarrow$ ".

perovskite being the dominant phase,  $La_2O_3$  being the second phase and  $La(OH)_3$  being the third phase. The main diffraction peaks of the  $La_2O_3$  phase and  $La(OH)_3$  phase are marked by the symbol " $\downarrow$ ". The diffraction peak positions of the three phases, according to the Rietveld powder diffraction profile fitting software (see below), are shown in Fig. 2. The upper series belongs to the perovskite phase, the middle series to the  $La_2O_3$  phase and the lower series to the

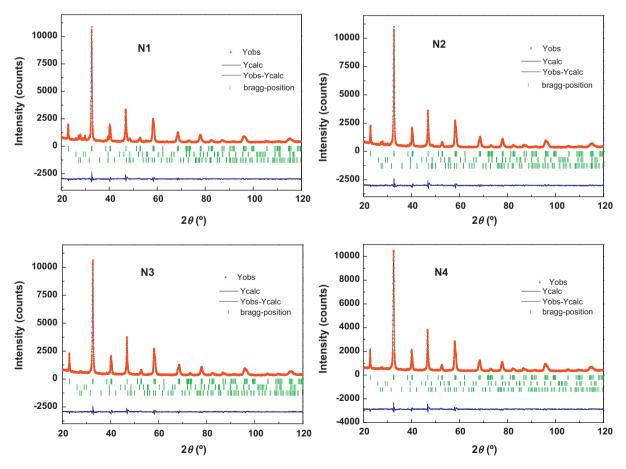
 $La(OH)_3$  phase. The reason for the appearance of the impurities,  $La_2O_3$  and  $La(OH)_3$ , is that the melting point of  $MnO_2$  is only 808 K, which is far lower than other starting materials (being higher than 1770 K), so that it is easy to escape from the raw materials during the samples preparation process.

By using the X'Pert HighScore Plus software to analyze the XRD patterns, the volume averaged crystallite diameters of the perovskite phase in the samples can be estimated. They were found to be 32, 38, 37, and 35 nm for samples N1, N2, N3, and N4, respectively. These values indicate that these samples were nanocrystalline composite materials.

Quantitative analyses of the three phases were performed using the main diffraction peak intensity ratios obtained using the X'Pert HighScore Plus software. We use the symbols "u" and "v" to designate the mass ratios of the La<sub>2</sub>O<sub>3</sub> and La(OH)<sub>3</sub> phases to the perovskite phase, respectively. According to X-ray diffraction theory we then have

$$u = \frac{W_{LO}}{W_P} = \frac{I_{LO}/K_{LO}}{I_P/K_P} = \frac{5.44}{8.20} \frac{I_{LO}}{I_P}, \quad v = \frac{W_{LH}}{W_P} = \frac{I_{LH}/K_{LH}}{I_P/K_P} = \frac{5.44}{6.86} \frac{I_{LH}}{I_P},$$
(1)

where  $W_{\rm LO}$ ,  $W_{\rm LH}$  and  $W_{\rm P}$  are the masses of the La<sub>2</sub>O<sub>3</sub> phase, La(OH)<sub>3</sub> phase and the perovskite phase respectively;  $I_{\rm LO}$ ,  $I_{\rm LH}$  and  $I_{\rm P}$  are the main diffraction peak intensities of the La<sub>2</sub>O<sub>3</sub>, La(OH)<sub>3</sub> and perovskite phases, respectively;  $K_{\rm LO}$ ,  $K_{\rm LH}$  and  $K_{\rm P}$ , with the values 8.20, 6.86 and 5.44, respectively, come from the *RIR* values for La<sub>2</sub>O<sub>3</sub>, La(OH)<sub>3</sub> and perovskite La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> on the ICDD standard cards 01-074-1144, 01-083-2034 and 01-089-4461. They each represent



**Fig. 2.** The Rietveld fitted results for the XRD diffraction data of the samples with the nominal composition  $La_{0.7}Sr_{0.3-x}MnO_{3-\delta}$  (x = 0.00, 0.05, 0.075, and 0.10, corresponding to samples N1, N2, N3, and N4). Here "•" represents the experimental intensity  $Y_{\text{obs}}$ . The bottom line shows the difference  $Y_{\text{obs}} - Y_{\text{cal}}$  between the experimental and fitted curve. The symbol "|" represents the diffraction peak positions of the three phases. The upper, middle and lower series belong to the perovskite.  $La_2O_3$  and  $La(OH)_3$  phases, respectively.

Table 1
The XRD main peak intensities  $I_P$ ,  $I_{LO}$  and  $I_{LH}$  of the perovskite phase,  $La_2O_3$  and  $La(OH)_3$  phases in the samples with a nominal composition  $La_{0.7}Sr_{0.3-x}MnO_{3-\delta}$ . The parameters  $c_{LO}$ ,  $c_{LH}$ , y, z and molar content of the perovskite phase M.

No.	X	$I_{\mathrm{P}}$	$I_{ m LO}$	$I_{ m LH}$	$c_{ m LO}$	$c_{ m LH}$	y	Z	М
N1	0.000	16276.83	569.30	923.48	5.23567	10.15191	0.02961	0.04924	0.92115
N2	0.050	16472.25	196.92	490.63	1.79383	5.34235	0.01007	0.02572	0.91421
N3	0.075	16011.43	52.59	243.06	0.49348	2.72621	0.00275	0.01305	0.90919
N4	0.100	16041.10	93.62	166.32	0.87365	1.85520	0.00475	0.00866	0.88658

the ratio of the main diffraction peak intensity of the compound in question to that of  $\alpha$ -Al<sub>2</sub>O<sub>3</sub> when the mass ratio is 1:1. Because there are different values of  $K_P$  for the perovskite phase as it occurs in the samples and for La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub>, the value of  $K_P$  used in this expression should be taken as inversely proportional to the density of the samples. Eq. (1) can then be rewritten as

$$\begin{split} u &= \frac{W_{\text{LO}}}{W_{\text{P}}} = 0.10037 \frac{I_{\text{LO}}}{I_{\text{P}}} \frac{5.44}{8.20} \frac{d_0 V_{\text{C}}}{m_{\text{P}}}, \quad v &= \frac{W_{\text{LH}}}{W_{\text{P}}} \\ &= 0.10037 \frac{I_{\text{LH}}}{I_{\text{P}}} \frac{5.44}{6.86} \frac{d_0 V_{\text{C}}}{m_{\text{P}}}, \end{split} \tag{2}$$

where  $m_P$  and  $V_C$  are the molar mass and the unit cell volume of the perovskite phase in our samples, expressed in units of grams (g) and Å<sup>3</sup>, respectively. Finally,  $d_0$  is the density of the standard La<sub>0.7</sub>Sr<sub>0.3</sub>MnO<sub>3</sub> material. Setting

$$c_{\text{LO}} = 0.10037 \frac{5.44}{8.20} \frac{I_{\text{LO}}}{I_{\text{P}}} d_0 V_{\text{C}}, \quad c_{\text{LH}} = 0.10037 \frac{5.44}{6.86} \frac{I_{\text{LH}}}{I_{\text{P}}} d_0 V_{\text{C}}, \quad (3)$$

then

$$\frac{W_{\text{LO}}}{W_{\text{P}}} = \frac{c_{\text{LO}}}{m_{\text{P}}}$$
 and  $\frac{W_{\text{LH}}}{W_{\text{P}}} = \frac{c_{\text{LH}}}{m_{\text{P}}}$ . (4)

We next designate the molar fractions of  $La_2O_3$  as y/2 and  $La(OH)_3$  as z. We assume that the missing molar content of  $MnO_2$  is equal to the sum of molar content of  $La_2O_3$  and  $La(OH)_3$  plus the doping level x.

Due to the samples were prepared using a solid-state reaction method with the highest heat treatment temperature being 1173 K, we assume that there were no vacancies in the perovskite phase. Therefore the molar content of Mn ions in the perovskite phase is M=1-x-y-z. The molar content of the perovskite phase is then (0.7-y-z+0.3-x+1-x-z-y)/2=1-x-y-z, and the oxygen molar content is 3(1-x-y-z). Note that the molar content of the Mn ions is equal to the molar content of the perovskite phase. The perovskite phase in the samples can therefore be written as  $\text{La}_{0.7-y-z}\text{Sr}_{0.3-x}\text{Mn}_{1-\delta/3}\text{O}_{3-\delta}$ , with  $\delta=3(x+y+z)$ . Following our earlier work [12,13], we have the various ion contents

$$L = 0.7 - y - z$$
,  $S = 0.3 - x$ ,  $M_3 = 0.7 - y - z$ ,  $M_4$   
=  $0.3 - x$ ,  $O_2 = 3(1 - x - y - z)$ , (5)

where L, S,  $M_3$ ,  $M_4$ , and  $O_2$  represent the molar contents of  $La^{3+}$ ,  $Sr^{2+}$ ,  $Mn^{3+}$ ,  $Mn^{4+}$  and  $O^{2+}$  ions, respectively. Therefore, we can represent the perovskite phase in our samples as  $La_{0.7-y-z}^{3+}Sr_{0.3-x}^{2+}Mn_{0.7-y-z}^{3+}Mn_{0.3-x}^{4+}O_{3(1-x-y-z)}^{2-}$ .

**Table 2** The molar contents L, S,  $M_3$ ,  $M_4$ , and  $O_2$  of  $La^{3+}$ ,  $Sr^{2+}$ ,  $Mn^{3+}$ ,  $Mn^{4+}$  and  $O^{2-}$  ions of the perovskite phase in the samples, and the  $Mn^{4+}$  ion content ratio,  $R_{M4}$ , at the B site.

L	S	$M_3$	$M_4$	$O_2$	$R_{\rm M4}$
0.62115	0.300	0.62115	0.300	2.76346	0.32568
0.66421	0.250	0.66421	0.250	2.74264	0.27346
0.68419	0.225	0.68419	0.225	2.72758	0.24747
0.68658	0.200	0.68658	0.200	2.65975	0.22558
	0.66421 0.68419	0.66421 0.250 0.68419 0.225	0.62115         0.300         0.62115           0.66421         0.250         0.66421           0.68419         0.225         0.68419	0.62115         0.300         0.62115         0.300           0.66421         0.250         0.66421         0.250           0.68419         0.225         0.68419         0.225	0.62115         0.300         0.62115         0.300         2.76346           0.66421         0.250         0.66421         0.250         2.74264           0.68419         0.225         0.68419         0.225         2.72758

The mass ratios of the  $La_2O_3$  phase and  $La(OH)_3$  phase to the perovskite phase are

$$\frac{W_{\rm LO}}{W_{\rm P}} = \frac{(y/2)m_{\rm LO}}{(1-x-y-z)m_{\rm P}}, \quad \frac{W_{\rm LH}}{W_{\rm P}} = \frac{zm_{\rm LH}}{(1-x-y-z)m_{\rm P}}, \tag{6}$$

where the molar masses of the La<sub>2</sub>O<sub>3</sub> phase and the La(OH)<sub>3</sub> phase are,  $m_{LO} = 325.809$  g and  $m_{LH} = 189.906$  g. From above Eqs. (4) and (6), we then have,

$$y = \frac{(1-x) \times 189.906c_{LO}}{30936.541977 + 189.906c_{LO} + 162.9045c_{LH}}, \quad z$$
$$= \frac{(1-x-y)c_{LH}}{189.906 + c_{LH}}. \tag{7}$$

Using Eqs. (3), (7) and the main peak intensity ratios of the three phases, we can obtain the parameter y and z. Table 1 lists the main peak intensities  $I_P$ ,  $I_{LO}$ , and  $I_{LH}$  of the perovskite, the La<sub>2</sub>O<sub>3</sub> and the La(OH)<sub>3</sub> phases, the parameters  $c_{LO}$ ,  $c_{LH}$ , y, z and the molar content M of the perovskite phase in the samples. Note again that M is also equal to the molar content of Mn ions in the perovskite phase as shown above.

The contents of the various ions as calculated using Eqs. (5) and (7), are shown in Table 2. In Table 2,  $R_{\rm M4}$  which represents the Mn<sup>4+</sup> ion content ratio at the B site of the ABO<sub>3</sub> structure, is defined by

$$R_{\rm M4} = \frac{M_4}{(M_3 + M_4)}. (8)$$

For convenience of interpretation, we also show in Table 3, the same ion contents normalized so that the  $O^{2-}$  content is 3, as is appropriate for the standard ABO<sub>3</sub> structure  $(La_{0.7-y-z}Sr_{0.3-x})_{1/(1-x-y-z)}MnO_3$ .

Using directly the contents of the various ions in Table 3, the XRD diffraction data were fitted using the Rietveld powder diffraction profile fitting technique [14]. In the fitting process, the contents of the various ions were not changed. The three equivalent positions (6a, 6b, and 18e) in the rhombohedral unit cell are occupied by 6 (La<sup>3+</sup>, Sr<sup>2+</sup>), 6 (Mn<sup>3+</sup>, Mn<sup>4+</sup>), and 18 O<sup>2-</sup>, respectively. The fitted results are shown in Fig. 2 and Table 4. It may be seen that the

Table 3
The molar contents L, S,  $M_3$ ,  $M_4$ , and  $O_2$  of La<sup>3+</sup>,  $Sr^{2+}$ ,  $Mn^{3+}$ ,  $Mn^{4+}$  and  $O^2$ —ions, the  $Mn^{4+}$  ion content ratio  $R_{M4}$ , at the B site, and molar mass  $m_P$  of the perovskite phase in the samples when the contents of the  $O^2$ —ions have been normalized to 3. The experimental moments  $σ_{SE}$  at 10 K under 2 T applied magnetic field, and theoritical moments  $σ_{SC}$  by the  $Mn^{3+}$  and  $Mn^{4+}$  content per molar perovskite phase in the samples.

No.	L	S	<i>M</i> <sub>3</sub>	$M_4$	02	R <sub>M4</sub>	$m_{ m P}$	$\sigma_{ ext{SE}}\left(\mu_{ ext{B}} ight)$	$\sigma_{ ext{SC}}\left(\mu_{ ext{B}} ight)$
N1	0.67432	0.32568	0.67432	0.32568	3	0.32568	225.14	3.47160	3.67432
N2	0.72654	0.27346	0.72654	0.27346	3	0.27346	227.82	3.48227	3.72654
N3	0.75253	0.24747	0.75253	0.24747	3	0.24747	229.15	3.48476	3.75253
N4	0.77442	0.22558	0.77442	0.22558	3	0.22558	230.27	3.51968	3.77442

**Table 4** The refined parameters for the rhombohedral perovskite phase in samples with the nominal composition  $La_{0.7}Sr_{0.3-x}MnO_{3-\delta}$ , the atomic coordinate at A, B and O site, calculated Mn–O bond lengths d and Mn–O–Mn bond angles  $\Theta$ , the profile factor  $R_{\rm p}$ , the weighted profile factor  $R_{\rm wp}$ , goodness of fit indicator s. Where S.G. is the space group, Z is the formular number per rhombohedral unit cell, a and c are lattice parameters, and V is the cell volume.

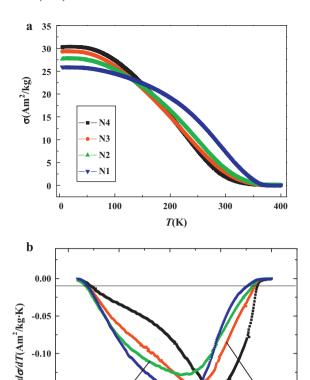
Samples	N1	N2	N3	N4
S.G.	R3c	R3̄c	R3̄c	R3c
Z	6	6	6	6
a (Å)	5.492(6)	5.499(0)	5.501(9)	5.506(5)
c (Å)	13.346(2)	13.347(7)	13.350(2)	13.362(2)
$V(Å^3)$	348.7(0)	349.5(4)	349.9(8)	350.8(8)
$A(La^{3+}Sr^{2+})$				
X	0	0	0	0
y	0	0	0	0
Z	0.25	0.25	0.25	0.25
$B(Mn^{3+}Mn^{4+})$				
X	0	0	0	0
У	0	0	0	0
Z	0	0	0	0
0				
X	0.468(3)	0.464(0)	0.462(2)	0.461(9)
У	0	0	0	0
Z	0.25	0.25	0.25	0.25
d (Å)	1.944(5)	1.948(3)	1.950(2)	1.952(5)
$\Theta$ ( $^{\circ}$ )	169.7(5)	168.3(5)	167.7(7)	167.4(0)
R <sub>p</sub> (%)	3.78	3.80	3.89	4.15
R <sub>wp</sub> (%)	4.79	4.90	4.92	5.25
S	1.16	1.19	1.18	1.24

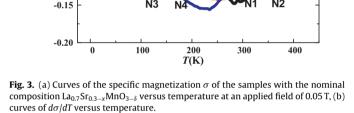
final fitting parameters, profile factor  $R_{\rm p}$ , weighted profile factor  $R_{\rm wp}$  and goodness of fit indicator s, are satisfactory, indicating that the contents of the various ions calculated by the above method are reasonable. From Table 4, it may be seen that the lattice parameters a and c, the cell volume V, Mn–O bond length d, the Mn–O–Mn bond angle  $\Theta$ , have all changed significantly.

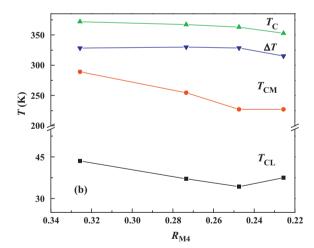
#### 3.2. Magnetic property analyses

Fig. 3(a) shows curves of the specific magnetization,  $\sigma$ , versus temperature at an applied field of 0.05 T. The curves show that with decreasing temperature, all samples exhibit a paramagnetic to ferromagnetic transition. Fig. 3(b) shows the curves  $d\sigma/dT$  versus temperature. The absolute values of  $d\sigma/dT$  begin to increase from near zero when  $T = T_{CL}$  (taken at  $d\sigma/dT = -0.1$ ), reach a maximum value when  $T = T_{CM}$ , and decrease to near zero when  $T = T_{C}$  (taken at  $d\sigma/dT = -0.1$ ). With increasing doping level x, the Curie temperature  $T_C$  shifts from 372 K for x = 0–353 K for x = 0.1. The transition region from ferromagnetism to paramagnetism,  $\Delta T = T_C - T_{CL}$ , is wider than 315 K. The main reason for this wide transition temperature range is that the volume averaged crystallite diameters of the perovskite phase are 32, 38, 37, and 35 nm in samples N1, N2, N3, and N4, respectively. The magnetic ordered temperature of the ion layers in the crystallites increases with decreasing the layer radius, with the result that the average angle between the Mn ionic moments decreases with decreasing the layer radius. We plot in Fig. 4 the relations of the magnetic ordered temperatures  $T_{\rm CL}$ ,  $T_{\rm CM}$ ,  $T_{\rm C}$ , and  $\Delta T$  versus the Mn<sup>4+</sup> ion content ratio,  $R_{\rm M4}$ , at the B site. From Fig. 4 and Table 4, it can be seen that the tendency of the change of the Curie temperature  $T_C$  is the same as that of the Mn–O–Mn bond angles  $\Theta$ , and is contrary to the Mn–O bond length d and unit cell volume  $V_C$ . These experimental phenomena can be explained using Zener's double exchange interaction between Mn<sup>3+</sup> and Mn<sup>4+</sup> ions [11,1,2].

Curves of the specific magnetization  $\sigma$  versus the applied field for the samples, tested at 300 K and 10 K, are shown in Fig. 5(a) and (b). It can be seen that the coercivity is very small in all our samples. At 10 K, the specific saturation magnetization  $\sigma_{SE}$  of the samples

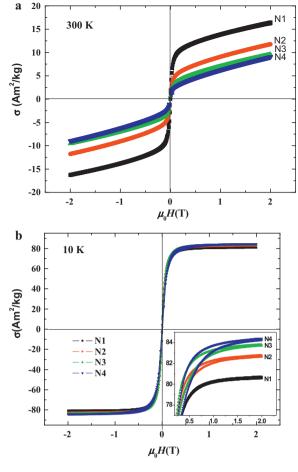






**Fig. 4.** The  $R_{\rm M4}$  dependences of the magnetic ordered temperatures  $T_{\rm CL}$ ,  $T_{\rm CM}$ ,  $T_{\rm C}$ , and  $\Delta T$ 

increases with decreasing Mn<sup>4+</sup> ion content ratio  $R_{\rm M4}$ , which is in agreement with the theoretical values,  $\sigma_{\rm SC}$ , as shown in Table 3. The value of  $\sigma_{\rm SE}$  is about 94% of  $\sigma_{\rm SC}$ , because the volume averaged crystallite diameters of the perovskite phase are between only 32 and 38 nm. At 300 K, the specific magnetization  $\sigma$  does not reach saturation until the applied field is increased to 2 T. The  $\sigma(2\,{\rm T})$  is trend as a function of  $R_{\rm M4}$  is contrary to that of  $\sigma_{\rm SE}$  at 10 K, resulting from the fact that 300 K is obviously higher than the magnetic ordered temperature  $T_{\rm CM}$  of the main phase in the perovskite nanocrystallites.



**Fig. 5.** Curves of the specific magnetization  $\sigma$  of the samples versus the applied field tested at 300 K (a) and 10 K (b).

# 3.3. Free energy considerations concerning the dependence of the unit cell volume in the perovskite phase on the doping level

From Table 4, it can be seen that the unit cell volume of the perovskite phase in the samples increases with increasing doping level x. This can be explained using the free energy of the unit cell [15,16]. In order to make the discussion of the cell volume more convenient, we transform the rhombohedra perovskite unit cell to the equivalent cubic cell, with lattice constant  $a_0$  (volume  $V_0=a_0^3$ ), in which there is only one ABO $_3$  group of atoms. Therefore, the problem of the doping level x dependence of the unit cell volume can be considered in terms of the dependence of the crystal lattice constant  $a_0$  on the doping level. The average free energy per equivalent cubic cell of the perovskite phase can be represented [15] as a sum,

$$F = U - TS = u_{\mathsf{C}} + u_{\mathsf{PR}} + u_{\mathsf{TV}},\tag{9}$$

$$\begin{split} u_{\mathsf{C}} &= u_{\alpha} + u_{\mathsf{b}} + u_{\mathsf{d}} = -\frac{q^2}{4\pi\varepsilon_0 a} \left[ \alpha + c_{\mathsf{b}} y \cos\left(\frac{\pi T^5}{2T_{\mathsf{C}}^5}\right) \right. \\ &\left. + \frac{c_{\mathsf{d}}}{a^2} e^{-E_{\mathsf{u}}/k_{\mathsf{B}}T} \right], \quad (c_{\mathsf{b}} = 0, \quad \text{when} \quad T > T_{\mathsf{C}}), \end{split} \tag{10}$$

$$u_{\rm PR} = \frac{\alpha q^2}{4\pi\varepsilon_0} \left( \frac{c_{\rm BO}}{(a/2)^i} + \frac{c_{\rm AO}}{(a/\sqrt{2})^j} + \frac{c_{\rm OO}}{(a/\sqrt{2})^m} + \frac{c_{\rm AB}}{(\sqrt{3}a/2)^n} \right), \quad (11)$$

$$u_{\text{TV}} = \left(\frac{a}{0.385}\right)^{5/3} \times 9Nk_{\text{B}}T\left(\frac{T}{\theta_{\text{D}}}\right)^{3} \int_{0}^{\theta_{\text{D}}/T} x^{2} \ln(1 - e^{-x}) dx.$$
 (12)

#### Table 5

Parameters used in free energy fitting the doping level dependence of the equivalent cubic cell lattice constant  $a_0$  for the perovskite phase in the samples. The parameters  $c_d$  and  $E_u$  are the thermal activation energy coefficient of the small polarons, and the thermal activation energy for small polarons, respectively. The Curie temperature  $T_C$ , Debye temperature  $\theta_D$ , the parameters  $c_b$ ,  $c_r$  and Madelung constant  $\alpha$  are also listed in the table

No.	х	<i>T</i> <sub>C</sub> (K)	$\theta_{\mathrm{D}}\left(\mathrm{K}\right)$	α	$c_{\rm b}$	$c_{\rm r}$	$c_{\mathrm{d}}$	E <sub>u</sub> (eV)
N1	0.00	372	437	45.67	0.3	0.9934	1	0.116
N2	0.05	367	427	45.56	0.3	0.9910	1	0.106
N3	0.075	363	420	45.51	0.3	0.9898	1	0.102
N4	0.10	354	413	45.48	0.3	0.9891	1	0.100

#### Table 6

The calculated lattice constant  $a_0$  and energies (eV) corresponding to the minimum average free energy  $F_0$  per equivalent cubic cell for the perovskite phase in the samples. Here  $u_{\alpha 0}$  is the ionic Coulomb energy,  $u_{b0}$  and  $u_{d0}$  are the additional energies produced by DE carriers and the thermal activation small polarons, respectively;  $u_{PRO}$  is the Pauli repulsion energy, and  $u_{TVO}$  is free energy corresponding to the thermal vibration energy of the crystal lattice.

No.	$a_0$	$F_0$	$u_{\alpha 0}$	$u_{b0}$	$u_{ m d0}$	$u_{\rm PRO}$	$u_{TV0}$
N1	0.38735	-23.0497	-27.1638	-0.04997	-0.04539	4.38619	-0.17664
N2	0.38767	-22.9825	-27.0779	-0.04096	-0.06512	4.38385	-0.18231
N3	0.38784	-22.9534	-27.0384	-0.03626	-0.07532	4.38282	-0.18618
N4	0.38815	-22.9198	-26.9960	-0.03017	-0.08227	4.37975	-0.19016

We can obtain the lattice constant  $a_0$  by minimizing the free energy F with respect to the lattice parameter a. In some ionic semiconductors, the effective unit charge q of the ions is far smaller than the electron charge e. Here we let q = 0.4e [16]. In Eq. (10),  $u_{\alpha}$  is the ionic Coulomb energy;  $u_{\rm b}$  is the additional energy produced by DE carriers, parameter y represents the number of Mn<sup>3+</sup> and  $Mn^{4+}$  ion pairs;  $u_d$  is the additional energy produced by small polarons. In Eq. (11),  $u_{PR}$  is the Pauli repulsion energy, in which the four terms are the repulsion energies in the O-B bonds, A-O bonds, O-O bonds and A-B bonds, respectively. The parameters  $c_{AO}$  =  $c_{OO}$  = 1.00 × 10<sup>-6</sup>, and the parameters  $c_{BO}$  = 3.30 × 10<sup>-5</sup>,  $c_{AB}$  = 1.10 × 10<sup>-5</sup>. The different denominators in the four terms represent the different bond lengths. The indices i, j, m, and n in Eq. (11) are related to the parameters  $r_k$ . Here  $r_k$  is the weighted average for the O and A or B ions,  $r_i = r(Mn^{3+}, O^{2-})M_3 + r(Mn^{4+}, O^{2-})M_4$ ,  $r_i = r(\text{La}^{3+}, \text{O}^{2-})L + r(\text{Sr}^{2+}, \text{O}^{2-})S$ ,  $r_m = r(\text{O}^{2-}, \text{O}^{2-})$ ,  $r_n = r_i + r_i - r_m$ . Here,  $r(Mn^{3+},O^{2-})$  is the sum of the Mn<sup>3+</sup> and O<sup>2-</sup> effective radii, and similarly for the other parameters  $r(\text{ion, O}^{2-})$ . The cation contents L, S,  $M_3$  and  $M_4$  come from Table 3. The indices are defined as

$$k = c_{\rm r}(-1.138 + 3.459r_k), \quad (k = i, m).$$
 (13)

$$k = -1.138 + 3.459r_k, \quad (k = j, m).$$
 (14)

The parameter  $c_r$  represents different repulsion power ratios between the shorter  $r_k$  ( $r_i$  or  $r_n$  relating to the Mn ions) and the longer  $r_k$  ( $r_j$  or  $r_m$  relating to non-Mn ions).

In Eq. (12),  $u_{\text{TV}}$  is the free energy corresponding to thermal vibration of the crystal lattice,  $\theta_{\text{D}}$  is Debye temperature of the perovskite phase, N=5, representing 5 atoms in every equivalent cubic cell. The parameters  $c_{\text{d}}$  and  $E_{\text{u}}$  are the thermal activation energy coefficient and the thermal activation energy, respectively, for small polarons. The Curie temperature,  $T_{\text{C}}$ , Debye temperature,  $\theta_{\text{D}}$  [1], Madelung constant,  $\alpha$  [16] and the parameters  $c_{\text{h}}$  [15] are listed in the Table 5.

Table 6 gives the calculated results: the lattice constant  $a_0$  (nm) of the equivalent cubic cell and the energies (eV) corresponding to the minimum average free energy  $F_0$  per equivalent cubic cell for the perovskite phase in the samples. Here  $u_{\alpha 0}$  is the ionic Coulomb energy,  $u_{b0}$  and  $u_{d0}$  are the additional energies produced by DE carriers and the small polarons by thermal activation, respectively;  $u_{PR0}$  is the Pauli repulsion energy, and  $u_{TV0}$  is free energy corresponding to the thermal vibration energy of the crystal lattice.

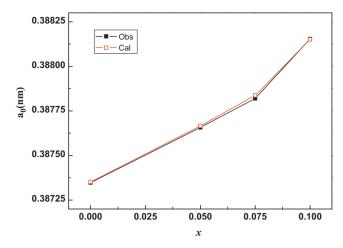


Fig. 6. The dependences of the calculated and experimental equivalent cubic cell lattice constant  $a_0$  in the perovskite manganites on self-doping level x.

Fig. 6 shows the calculated and experimental dependences of the equivalent cubic cell lattice constant  $a_0$  on the Mn<sup>4+</sup> ion content ratio  $R_{M4}$  of the perovskite phase in the samples. It can be seen that the experimental values and the calculated values match very well.

#### 4. Conclusions

We have prepared manganites with the nominal composition  $La_{0.7}Sr_{0.3-x}MnO_{3-\delta}$  ( $0 \le x \le 0.10$ ) by a solid-state reaction method with the highest heat treatment temperature being 1173 K. Using XRD analysis, the samples were verified as being three phase composities  $La_{0.7-y-z}Sr_{0.3-x}Mn_{1-\delta/3}O_{3-\delta}/(La_2O_3)_{y/2}/(La(OH)_3)_z$ , with  $\delta = 3(x+y+z)$ . We have calculated the ion content ratios at the A, B and O sites of the perovskite phase, and the ratios were shown to be reasonable through Rietveld fitting, analysis of the magnetic property. It was found that the dependences of the Curie temperature  $T_{\rm C}$  on the Mn<sup>4+</sup> ion content  $R_{\rm M4}$  at the B sites, are similar to those of the typical perovskite  $La_{1-x}Sr_xMnO_3$ . At 10 K, the specific saturation magnetization  $\sigma_{\rm SE}$  of the samples increases with decreasing  $Mn^{4+}$  ion content ratio  $R_{M4}$ , which is in accord with changes in the theoretical value  $\sigma_{SC}$ . The volume averaged crystallite diameters of the perovskite phase in the samples are between 32 and 38 nm, resulting in the observation that all samples display a very wide paramagnetic-ferromagnetic transition region with decreasing temperature. We calculated the free energy of the perovskite phase, and explained successfully the dependence of the unit cell volume of the perovskite phase on the doping level.

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#### References

- [1] M.B. Salamon, M. Jaime, Rev. Mod. Phys. 73 (2001) 583.
- E. Dagotto, T.i. Hotta, A. Moreo, Phys. Rep. 344 (2001) 1-153.
- [3] R. Mahendiran, S.K. Tiwary, A.K. Raychaudhuri, T.V. Ramakrishnan, R. Mahesh, N. Rangavittal, C.N.R. Rao, Phys. Rev. B 53 (1996) 3348-3358.
- [4] G.D. Tang, D.L. Hou, W. Chen, X. Zhao, W.H. Qi, Appl. Phys. Lett. 90 (2007) 144101.
- B. Lin, S. Wang, H. Liu, K. Xie, H. Ding, M. Liu, G. Meng, J. Alloys Compd. 472 (2009) 556-558.
- L. Bi, S. Zhang, B. Lin, S. Fang, C. Xia, W. Liu, J. Alloys Compd. 473 (2009) 48-52.
- S. Wang, B. Lin, Y. Chen, X. Liu, G. Meng, J. Alloys Compd. 479 (2009) 764-768.
- [8] M. Raama Varma, P. Nisha, P.C. Rajath Varma, J. Alloys Compd. 457 (2008)
- [9] J. Deng, J. Chen, R. Yu, G. Liu, X. Xing, J. Alloys Compd. 472 (2009) 502-506.
- [10] S.-F. Wang, Y.-R. Wang, Y.-C. Wu, Y.-J. Liu, J. Alloys Compd. 480 (2009) 499-504.
- [11] C. Zener, Phys. Rev. 82 (1951) 403.
- S.P. Liu, G.D. Tang, P. Hao, L.Q. Xu, Y.G. Zhang, W.H. Qi, X. Zhao, D.L. Hou, W. Chen, J. Appl. Phys. 105 (2009) 013905.
- G.D. Tang, D.L. Hou, W. Chen, P. Hao, G.H. Liu, S.P. Liu, X.L. Zhang, L.Q. Xu, Appl. Phys. Lett. 91 (2007) 152503.
- H.M. Rietveld, J. Appl. Crystallogr. 2 (1969) 65. G.D. Tang, S.P. Liu, W. Zhao, Y.G. Zhang, D.H. Ji, Y.F. Li, W.H. Qi, W. Chen, D.L. Hou, Appl. Phys. Lett. 95 (2009) 121906.
- G.D. Tang, D.L. Hou, Z.Z. Li, X. Zhao, W.H. Qi, S.P. Liu, F.W. Zhao, Appl. Phys. Lett. 89 (2006) 261919.